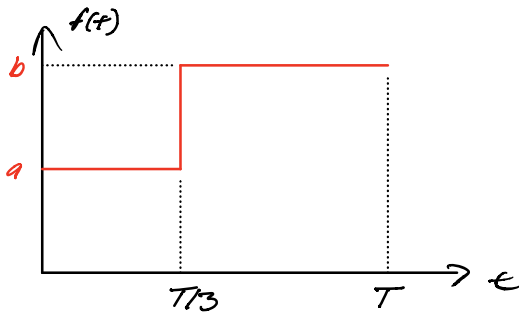


Problem 1



All calculations done in MATLAB and the expressions may differ from hand calculations due to the simplification

Approximate $f(t)$ as a Fourier Series

$$f(t) \cong f_0 + \sum_{n=1}^{\infty} f_{s,n} \sin n\Omega t + f_{c,n} \cos n\Omega t$$

$$f_0 = \frac{1}{T} \int_0^T f(t) dt = \frac{1}{T} \left\{ \int_0^{T/3} a dt + \int_{T/3}^T b dt \right\}$$

$$f_0 = \frac{1}{3} (a + 2b)$$

$$f_{s,n} = \frac{2}{T} \int_0^T f(t) \sin n\Omega t dt$$

$$f_{s,n} = \frac{2}{T} \left\{ \int_0^{T/3} a \sin n\Omega t dt + \int_{T/3}^T b \sin n\Omega t dt \right\}$$

$$f_{s,n} = \frac{1}{\pi} \left\{ 2a \sin\left(\frac{2\pi n}{3}\right) - 4b \cos\left(\frac{2\pi n}{3}\right) \left(\cos\left(\frac{2\pi n}{3}\right)^2 - 1\right) \right\}$$

$$f_{c,n} = \frac{2}{T} \int_0^T f(t) \cos n\Omega t dt$$

$$f_{c,n} = \frac{2}{T} \left\{ \int_0^{T/3} a \cos n\Omega t dt + \int_{T/3}^T b \cos n\Omega t dt \right\}$$

$$f_{c,n} = \frac{1}{\pi} \left\{ a \sin\left(\frac{2\pi n}{3}\right) + b \left(2 \sin\left(\frac{2\pi n}{3}\right) - 4 \sin\left(\frac{2\pi n}{3}\right)^3 \right) \right\}$$

Now, solving the equation

$$m\ddot{x} + c\dot{x} + kx = f(t) \quad x(0) = x_0 \quad \text{and} \quad \dot{x}(0) = v_0$$

using the convolution integral

The force can be written as a piecewise function

$$f(t) = \begin{cases} a & , 0 \leq t \leq T/3 \\ b & T/3 \leq t \leq T \end{cases}$$

from $0 \leq t \leq T/3$

$$x(t) = e^{-3\omega h t} \left(x_0 \cos \omega t + \frac{(v_0 + x_0 \omega h)}{\omega d} \sin \omega t \right) + \int_0^t \cancel{f(\tau)} \frac{e^{-3\omega h (t-\tau)}}{m \omega d} \sin \omega d (t-\tau) d\tau$$

$$x(t) = e^{-3\omega h t} \left(x_0 \cos \omega t + \frac{(v_0 + x_0 \omega h)}{\omega d} \sin \omega t \right) + \frac{a}{m \omega h^2} - \frac{a e^{-3\omega h t} (\omega d \cos \omega t + \omega h \sin \omega t)}{m \omega h^2}$$

from $0 \leq t < T$

$$x(t) = e^{-3\omega h t} \left(x_0 \cos \omega t + \frac{(v_0 + x_0 \omega h)}{\omega d} \sin \omega t \right) + \int_0^{T/3} \cancel{f(\tau)} \frac{e^{-3\omega h (t-\tau)}}{m \omega d} \sin \omega d (t-\tau) d\tau + \int_0^T \cancel{f(\tau)} \frac{e^{-3\omega h (t-\tau)}}{m \omega d} \sin \omega d (t-\tau) d\tau$$

$$\begin{aligned}
 x(t) = & e^{-\xi \omega_n t} \left(x_0 \cos \omega_d t + \frac{(v_0 + \xi \omega_n x_0)}{\omega_d} \sin \omega_d t \right) \\
 & + \frac{D_0}{\omega_n^2} + \frac{(a-b) \sigma_2 \sigma_4}{m \omega_d \omega_n^2} \\
 & - \frac{a c}{m \omega_d \omega_n^2} \left(\omega_d \cos \omega_d t + \xi \omega_n \sin \omega_d t \right)
 \end{aligned}$$

$$\sigma_2 = e^{-\xi \omega_n (t + T/3)}$$

$$\sigma_4 = \omega_d \cos \omega_d (T/3 - t) - \xi \omega_n \sin \omega_d (T/3 - t)$$

Now, solving the equation

$$m\ddot{x} + c\dot{x} + kx = f(t) \quad x(0) = x_0 \quad \text{and} \quad \dot{x}(0) = v_0$$

using a Fourier series

$$f(t) \approx f_0 + \sum_{n=1}^{\infty} \{ f_{sn} \sin n\Omega t + f_{cn} \cos n\Omega t \}$$

The homogeneous solution

$$x_h(t) = e^{-\zeta\omega_n t} (A_0 \cos \omega_d t + B_0 \sin \omega_d t)$$

$$x_p = C + \sum_{n=1}^{\infty} A_n \cos n\Omega t + \sum_{n=1}^{\infty} B_n \sin n\Omega t$$

We can take plug $x_p, \dot{x}_p, \ddot{x}_p$ into governing equation

$$\begin{aligned} & \sum_{n=1}^{\infty} \{ -n^2 A_n \Omega^2 \sin(n\Omega t) - n^2 B_n \Omega^2 \cos(n\Omega t) \\ & + 2\zeta\omega_n \Omega n A_n \cos(n\Omega t) - 2\zeta\omega_n \Omega n B_n \sin(n\Omega t) \\ & + \omega_n^2 (C + \omega_n^2 A_n \sin(n\Omega t) + \omega_n^2 B_n \cos n\Omega t) \\ & = f_0 + f_{sn} \sin n\Omega t + f_{cn} \cos n\Omega t \} \end{aligned}$$

Separate out terms

$$\text{constant:} \quad \omega_n^2 C = f_0$$

$$\sin n\Omega t: \quad \sum_{n=1}^{\infty} (-n\Omega^2 + \omega_n^2) A_n - 2\zeta\omega_n \Omega n B_n = f_{sn}$$

$$\cos n\Omega t: \quad \sum_{n=1}^{\infty} (-n\Omega^2 + \omega_n^2) B_n + 2\zeta\omega_n \Omega n A_n = f_{cn}$$

Solving $C = f_0 / \omega_0^2$

$$A_n = \frac{(\omega_n^2 - n^2 \Omega^2) f_{0n} + 2\zeta \omega_n n \Omega f_{0n}}{\Delta_n}$$

$$B_n = \frac{(\omega_n^2 - n^2 \Omega^2) f_{0n} - 2\zeta \omega_n n \Omega f_{0n}}{\Delta_n}$$

$$\Delta_n = n^2 \Omega^4 + \omega_n^4 + n^2 \Omega^2 \omega_n^2 (4\zeta^2 - 2)$$

Now the total solution is

$$x(t) = e^{-\zeta \omega_0 t} (A_0 \cos \omega_0 t + B_0 \sin \omega_0 t) + C + \sum_{n=1}^{\infty} A_n \cos n \Omega t + \sum_{n=1}^{\infty} B_n \sin n \Omega t$$

$$\dot{x}(t) = -\zeta \omega_0 e^{-\zeta \omega_0 t} (A_0 \cos \omega_0 t + B_0 \sin \omega_0 t) + e^{-\zeta \omega_0 t} (-\omega_0 A_0 \sin \omega_0 t + \omega_0 B_0 \cos \omega_0 t) + \sum_{n=1}^{\infty} -n \Omega A_n \sin n \Omega t + \sum_{n=1}^{\infty} n \Omega B_n \cos n \Omega t$$

Solving for A_0 and B_0 using $x(0) = x_0$ and $\dot{x}(0) = v_0$

$$x_0 = A_0 + \overset{f_0 / \omega_0^2}{C} + \sum_{n=1}^{\infty} A_n$$

$$v_0 = -\zeta \omega_0 (A_0) + \omega_0 B_0 + \sum_{n=1}^{\infty} n \Omega B_n$$

Solving we can write as a matrix equations of the form $Ax = b$

$$\begin{bmatrix} 1 & 0 \\ -\zeta \omega_0 & \omega_0 \end{bmatrix} \begin{Bmatrix} A_0 \\ B_0 \end{Bmatrix} = \begin{Bmatrix} -f_0 / \omega_0^2 - \sum_{n=1}^{\infty} A_n + x_0 \\ -\sum_{n=1}^{\infty} n \Omega B_n + v_0 \end{Bmatrix}$$

$$\begin{Bmatrix} A_0 \\ b_0 \end{Bmatrix} = \frac{1}{\omega_d} \begin{bmatrix} \omega_d & +3\omega_h \\ 0 & 1 \end{bmatrix} \begin{Bmatrix} x_0 - \frac{f_0}{\omega_h^2} - \sum_{n=1}^{\infty} A_n \\ v_0 - \sum_{n=1}^{\infty} n \omega_n B_n \end{Bmatrix}$$

In summary, the response can be written as follows
for convolution integral

$$0 \leq t \leq T/3$$

$$x(t) = e^{-3\omega_h t} \left(x_0 \cos \omega_d t + \frac{(v_0 + f_0 \omega_h)}{\omega_d} \sin \omega_d t \right) + \frac{a}{m \omega_h^2} - \frac{a e^{-3\omega_h t} (\omega_d \cos \omega_d t + \omega_h 3 \sin \omega_d t)}{m \omega_d \omega_h^2}$$

$$T/3 \leq t \leq T$$

$$x(t) = e^{-3\omega_h t} \left(x_0 \cos \omega_d t + \frac{(v_0 + f_0 \omega_h)}{\omega_d} \sin \omega_d t \right) + \frac{b_0}{\omega_h^2} + \frac{(a-b)\sigma_2 \sigma_4}{m \omega_d \omega_h^2} - \frac{a e^{-3\omega_h t}}{m \omega_d \omega_h^2} (\omega_d \cos \omega_d t + 3\omega_h \sin \omega_d t)$$

$$\sigma_2 = e^{-3\omega_h (t + T/3)}$$

$$\sigma_4 = \omega_d \cos \omega_d (T/3 - t) - 3\omega_h \sin \omega_d (T/3 - t)$$

and for the Fourier Series

$$v_n(t) = e^{-3\omega_n t} (A_0 \cos \omega_n t + B_0 \sin \omega_n t) +$$

$$C + \sum_{n=1}^{\infty} A_n \cos n\Omega t + \sum_{n=1}^{\infty} B_n \sin n\Omega t$$

where

$$C = f_0 / \omega_n^2$$

$$A_n = \frac{(\omega_n^2 - n\Omega^2) f_{cn} + 2\zeta \omega_n n\Omega f_{sn}}{\Delta_n}$$

$$B_n = \frac{(\omega_n^2 - n\Omega^2) f_{sn} - 2\zeta \omega_n n\Omega f_{cn}}{\Delta_n}$$

$$\Delta_n = n\Omega^4 + \omega_n^4 + n^2 \Omega^2 \omega_n^2 (4\zeta^2 - 2)$$

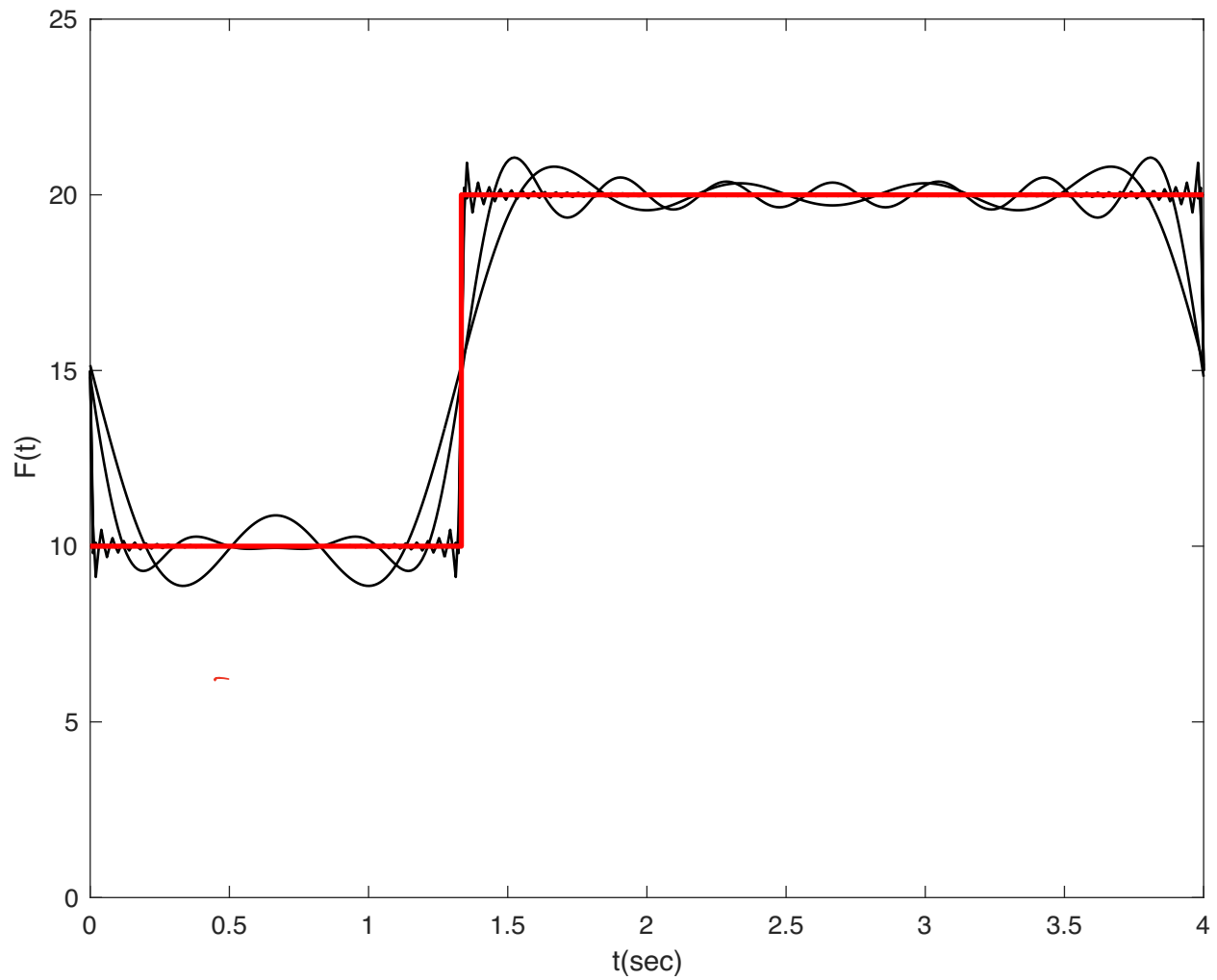
$$f_0 = \frac{1}{3} (a + 2b)$$

$$f_{cn} = \frac{1}{\pi} \left\{ 2a \sin\left(\frac{2\pi n}{3}\right)^2 - 4b \cos\left(\frac{2\pi n}{3}\right) (\cos\left(\frac{2\pi n}{3}\right)^2 - 1) \right\}$$

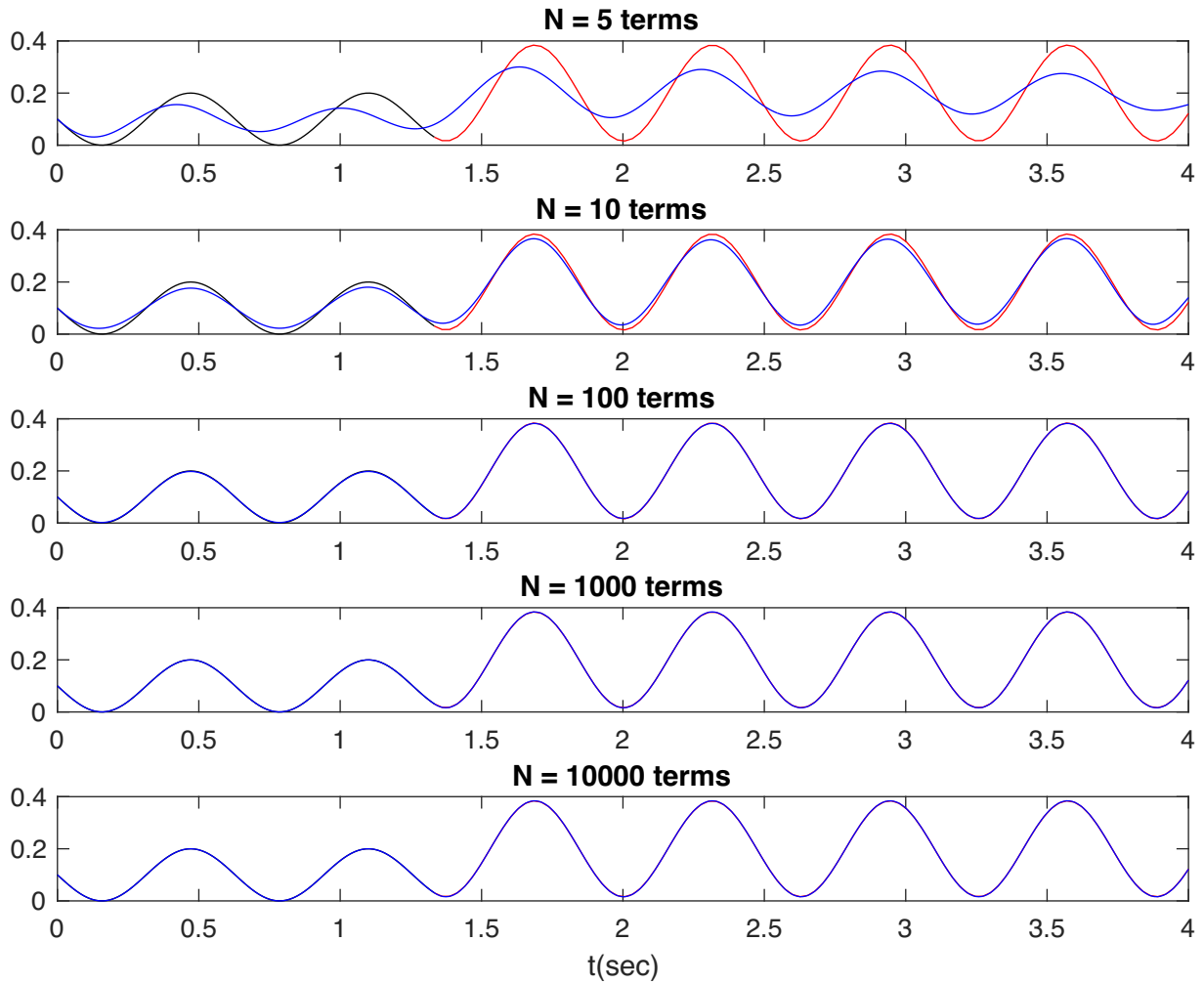
$$f_{sn} = \frac{1}{\pi} \left\{ a \sin\left(\frac{2\pi n}{3}\right) + b (2 \sin\left(\frac{2\pi n}{3}\right) - 4 \sin\left(\frac{2\pi n}{3}\right)^3) \right\}$$

$$\begin{Bmatrix} A_0 \\ B_0 \end{Bmatrix} = \frac{1}{\overline{\omega_n}} \begin{bmatrix} \omega_n & +3\omega_n \\ 0 & 1 \end{bmatrix} \begin{Bmatrix} f_0 - f_0/\omega_n^2 - \sum_{n=1}^{\infty} A_n \\ f_0 - \sum_{n=1}^{\infty} n\Omega B_n \end{Bmatrix}$$

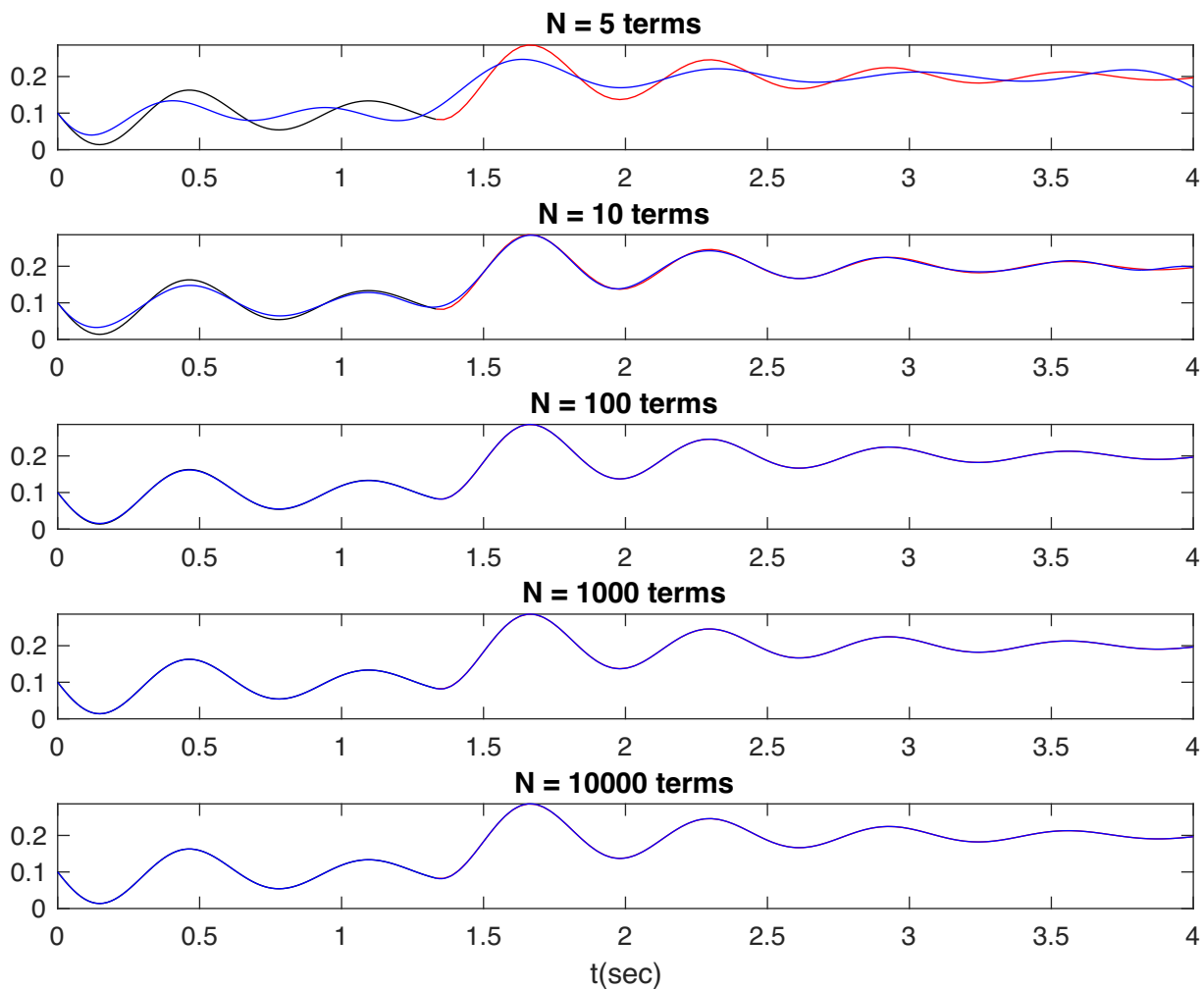
Fourier fit for $N=5, 10, 100, 1000, \text{ and } 10000$



response when $c=0$



Response when $c=2$



```

% Oscillator Parameters
close all
mass = 1;
damp = 2;
stiff = 100;

% Mass Normalization
Omega_n = sqrt(stiff/mass);
Zeta = damp/2*1/sqrt(mass*stiff);
Omega_d = Omega_n*sqrt(1-Zeta^2);

% Initial Conditions
X0 = 0.1;
V0 = -1.0;

% Force Parameters
a_actual = 10;
b_actual = 20;
T_actual = 4;
Omega_actual = 2*pi/T_actual;

% Define Force to Plot
F_actual = [a_actual a_actual b_actual b_actual];
t_actual = [0 T_actual/3 T_actual/3 T_actual];

figure(1);
line(t_actual, F_actual, 'color', 'r', 'linewidth',2)
ylabel('F(t)')
xlabel('t(s)')
box on

syms a b T t Omega real
syms n integer

%% Fourier Series %%%%%%%%%%%%%%%
% Constant Term
fo = 1/T*(int(a,t,[0,T/3])+ int(b, t,[T/3,T]));
fo = simplify(fo)
fof = matlabFunction(fo); % Change to a function
fo;
% Cosine Term
fc1 = 2/T*int(a*cos(n*2*pi/T*t), t,[0,T/3]);
fc2 = 2/T*int(b*cos(n*2*pi/T*t), t,[T/3,T]);
assume(n,'integer')
fc1 = simplify(fc1);
fc2 = simplify(fc2);
fc = fc1 + fc2;
fcf = matlabFunction(fc);
fc
% Sine Term
fxp1 = 2/T*int(a*sin(n*2*pi/T*t),t,[0,T/3]);
fxp2 = 2/T*int(b*sin(n*2*pi/T*t), t,[T/3,T]);
fxp1 = simplify(fxp1);
fxp2 = simplify(fxp2);
fs = fxp1 + fxp2;
fsf = matlabFunction(fs); % convert to inline function
fs

```

```

% Number of Fourier Terms
Nv = [5 10 100 1000 10000];
tv = linspace(0, T_actual, 400);

for i = 1:length(Nv)
    N = Nv(i);
    Fo = fof(a_actual, b_actual);
    Fs = 0;
    Fc = 0;
    for nv = 1:N
        Fs = Fs + fsf(a_actual, b_actual, nv)*sin(nv*2*pi*tv/T_actual);
        Fc = Fc + fcf(a_actual, b_actual, nv)*cos(nv*2*pi*tv/T_actual);
    end

    F_fourier = Fo+Fs+Fc;

    figure(2)
    line(tv, F_fourier, 'color', 'k', 'linewidth', 1)
    box on
    axis([0 4 0 25])
end
line(t_actual, F_actual, 'color', 'r', 'linewidth', 2)
xlabel('t(sec)')
ylabel('F(t)')

%% Convolution Integral
syms h m zeta omega_n omega_d tau a b x0 v0 reals
assume(zeta>0 & zeta<1)

h = exp(-zeta*omega_n*(t-tau))*1/(m*omega_d)*sin(omega_d*(t-tau));

xh = exp(-zeta*omega_n*t)*((zeta*omega_n*x0+v0)/(omega_d)*sin(omega_d*t)+x0*cos(
(omega_d*t))

x1 = int(a*h,tau,[0,t]);
x1 = simplify(x1+xh) %Add in homogenous solution
x1c = matlabFunction(x1) %Convert to Matlab Function

x2a = int(a*h,tau,[0,T/3]);
x2b = int(b*h,tau,[T/3,t]);
x2 = simplify(x2a+x2b+xh)
x2c = matlabFunction(x2) %Convert to Matlab Function

%% Harmonic Balance with Fourier Series
syms x(t) FFo FFsn FFcn C1 C2 An Bn C A0 B0 xo vo

% Homogenous Solution
xhf = exp(-zeta*omega_n*t)*(A0*cos(omega_d*t) + B0*sin(omega_d*t));
Dxhf = diff(xhf,t)

% Particular Soution
xpf = C+ An*sin(n*Omega*t) + Bn*cos(n*Omega*t)
Dxpf = diff(xpf,t)
Dx2xpf = diff(Dxpf,t)

Eqn1 = Dx2xpf + 2*zeta*omega_n*Dxpf +omega_n^2*xpf -(FFo + FFsn*sin(n*Omega*t) + FFcn*cos(

```

```

(n*Omega*t))

Sin_terms = coeffs(Eqn1, sin(n*Omega*t));
Sin_terms = Sin_terms(2)

Cos_terms = coeffs(Eqn1, cos(n*Omega*t));
Cos_terms = Cos_terms(2)

Constant_terms = subs(Eqn1, [sin(n*Omega*t) cos(n*Omega*t)], [0, 0])

Csol = solve(Constant_terms ==0,C)
Sol=solve(Sin_terms==0, Cos_terms==0,[An Bn])

ASol=Sol.An;
BSol=Sol.Bn;

%IC homogeneous
xho=subs(xhf,t,0)
vho=subs(Dxhf,t,0)

AAn = matlabFunction(ASol)
BBn = matlabFunction(BSol)

for ii = 1:length(Nv)

    N    = Nv(ii);
    Xpf  = 0; % The particular solution
    Xpfo = 0; % Particular solution
    Vpfo = 0;

    for nv = 1:N
        fsn = fsf(a_actual, b_actual, nv);
        fcn = fcf(a_actual, b_actual, nv);
        an  = AAn(fcn, fsn, Omega_actual, nv, Omega_n, Zeta);
        bn  = BBn(fcn, fsn, Omega_actual, nv, Omega_n, Zeta);

        Xpf = Xpf + an*sin(nv*Omega_actual*tv) + bn*cos(nv*Omega_actual*tv);

        Xpfo = Xpfo + bn;
        Vpfo = Vpfo + an*nv*Omega_actual;
    end
    fon = fof(a_actual, b_actual);
    Xpfo = Xpfo + fon/Omega_n^2;
    Xpf  = Xpf  + fon/Omega_n^2;

    F    = [X0-Xpfo;V0-Vpfo];
    IcM  = [1 0; -Omega_n*Zeta Omega_d];
    Icv  = IcM\F;

    Ao=Icv(1);
    Bo=Icv(2);

    Xf = exp(-Zeta*Omega_n*tv).*(Ao*cos(Omega_d*tv) + Bo*sin(Omega_d*tv));
    Xf = Xf + Xpf;

    %Time Vectors for Convolutions- notreally efficient to put in the loop
    %but it makes the code easier to read
    tv1 = linspace(0,T_actual/3,100);

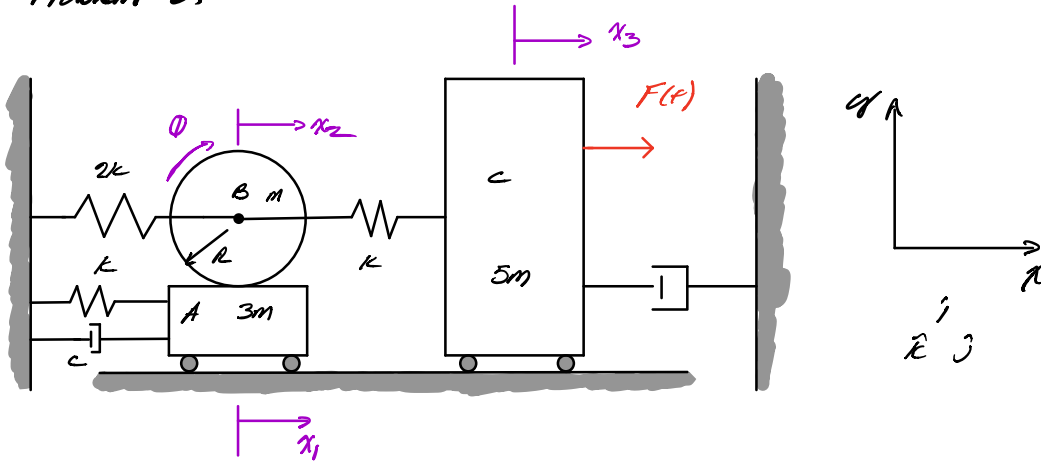
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tv2 = linspace(T_actual/3,T_actual,100);
X1c = x1c(a_actual, mass, Omega_d, Omega_n, tv1, V0, X0, Zeta);
X2c = x2c(T_actual, a_actual, b_actual, mass, Omega_d, Omega_n, tv2, V0, X0, Zeta);

figure(3)
subplot(length(Nv),1,ii)
title(['N = ',num2str(N), ' terms'])
line(tv1,X1c,'color','k') % Plot 0<t<T/3 Convolution
line(tv2,X2c,'color','r') % Plot T/3<t<T Convolution
line(tv,Xf,'color','b') % Plot Fourier Series
box on
end
xlabel('t(sec)')

box on
```

Problem 2



Begin by deriving the equations of motion using the linearized Lagrange's Methods

$$T = \frac{1}{2} (3m) \dot{x}_1^2 + \frac{1}{2} m \dot{x}_2^2 + \frac{1}{2} I_B \dot{\theta}^2 + \frac{1}{2} (5m) \dot{x}_3^2$$

$$I_B = \frac{1}{2} m R^2$$

Next by kinematics

$$\vec{v}_2 = \vec{v}_1 + \vec{\omega} \times \vec{r}_{e1}$$

where $\vec{v}_2 = \dot{x}_2 \hat{i}$, $\vec{\omega} = -\dot{\theta} \hat{k}$, $\vec{r}_{e1} = R \hat{j}$, and $\vec{v}_1 = \dot{x}_1 \hat{i}$

$$\dot{x}_2 \hat{i} = \dot{x}_1 \hat{i} + (-\dot{\theta} \hat{k}) \times R \hat{j} = (\dot{x}_1 + \dot{\theta} R) \hat{i}$$

$$\text{Now, } \underline{\dot{\theta} = (\dot{x}_2 - \dot{x}_1) / R}$$

$$T = \frac{1}{2} (3m) \dot{x}_1^2 + \frac{1}{2} m \dot{x}_2^2 + \frac{1}{2} \left(\frac{1}{2} m R^2 \right) \frac{(\dot{x}_2 - \dot{x}_1)^2}{R^2} + \frac{1}{2} (5m) \dot{x}_3^2$$

$$T = \frac{1}{2} (3m) \dot{x}_1^2 + \frac{1}{2} m \dot{x}_2^2 + \frac{1}{2} \left(\frac{1}{2} m \right) (\dot{x}_2 + \dot{x}_1 - 2\dot{x}_1 \dot{x}_2) + \frac{1}{2} (5m) \dot{x}_3^2$$

$$T = \frac{1}{2} (3m + \frac{1}{2}m) \dot{x}_1^2 + \frac{1}{2} (m + \frac{1}{2}m) \dot{x}_2^2 + \frac{1}{2} (-m) \dot{x}_1 \dot{x}_2 + \frac{1}{2} (5m) \dot{x}_3^2$$

$$T = \frac{1}{2} \underset{m_{11}}{(1/2 m)} \dot{x}_1^2 + \frac{1}{2} \underset{m_{12}}{(3/2 m)} \dot{x}_2^2 + \frac{1}{2} \underset{m_{12}+m_{21}}{(-m)} \dot{x}_1 \dot{x}_2 + \frac{1}{2} \underset{m_{33}}{(5m)} \dot{x}_3^2$$

The mass matrix can be written as

$$[m] = m \begin{bmatrix} 1/2 & -1/2 & 0 \\ -1/2 & 3/2 & 0 \\ 0 & 0 & 5 \end{bmatrix} = m \begin{bmatrix} \bar{m} \end{bmatrix}$$

$\sim \bar{m}$

Now, the potential energy can be written as

$$U = \frac{1}{2} k x_1^2 + \frac{1}{2} (2k) x_2^2 + \frac{1}{2} k (x_3 - x_2)^2$$

Using linearized Lagrange Method

$$k_{11} = \frac{\partial^2 U}{\partial x_1^2} = k$$

$$k_{12} = k_{21} = \frac{\partial^2 U}{\partial x_2 \partial x_1} = 0$$

$$k_{13} = k_{31} = \frac{\partial^2 U}{\partial x_3 \partial x_1} = 0$$

$$k_{22} = \frac{\partial^2 U}{\partial x_2^2} = 2k + k = 3k$$

$$k_{23} = k_{32} = \frac{\partial}{\partial x_2} (2k x_2 - k(x_3 - x_2)) = -k$$

$$k_{33} = \frac{\partial^2 U}{\partial x_3^2} = k$$

$$[k] = k \begin{bmatrix} 1 & 0 & 0 \\ 0 & +3 & -1 \\ 0 & -1 & 1 \end{bmatrix} = k \begin{bmatrix} \bar{k} \end{bmatrix}$$

$\sim \bar{k}$

The Rayleigh damping

$$R = \frac{1}{2} c \dot{x}_1^2 + \frac{1}{2} c \dot{x}_3^2$$

$$c_{11} = c, \quad c_{12} = c_{21} = 0, \quad c_{13} = c_{31} = 0$$

$$c_{22} = 0, \quad c_{23} = 0$$

$$c_{33} = c$$

$$[C] = c \begin{bmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 1 \end{bmatrix} = c [\bar{c}]$$

Virtual work

$$dW = \vec{F} \cdot d\vec{x}_3 \uparrow = \underbrace{F(x)}_{Q_3} \cdot dx_3$$

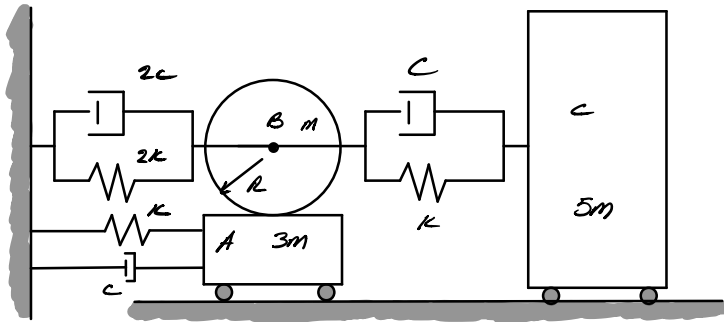
$$\vec{F} = \begin{Bmatrix} 0 \\ 0 \\ F \end{Bmatrix}$$

$$[m] \ddot{\vec{x}} + [c] \dot{\vec{x}} + [k] \vec{x} = \vec{F}$$

The system is not proportionally damped

Note the system cannot be made to be proportional to the mass matrix due to the rotational inertia which causes inertial coupling between x_1 and x_2

However, the damping can be made proportional to the stiffness



In determining the mode shapes set $[c] = 0$
 assume $\vec{x} = \vec{X} e^{i\omega t}$

$$[m] \ddot{\vec{x}} + [k] \vec{x} = \vec{0} \rightarrow \underbrace{[-\omega^2 [m] + [k]]}_0 \vec{X} = \vec{0}$$

$$\text{Det}(0) = -25 m^3 \omega^6 + 65 km^2 \omega^4 - 47/2 k^2 m \omega^2 + 2k^3 = 0$$

In MATLAB

$$\omega_1 = 0.359 \sqrt{k/m}$$

$$\omega_2 = 0.533 \sqrt{k/m}$$

$$\omega_3 = 1.474 \sqrt{k/m}$$

$$\vec{X}_1 = \frac{1}{\sqrt{m}} \begin{Bmatrix} -0.018 \\ 0.156 \\ 0.438 \end{Bmatrix} \quad \vec{X}_2 = \frac{1}{\sqrt{m}} \begin{Bmatrix} 0.531 \\ -0.015 \\ 0.036 \end{Bmatrix}$$

$$\vec{X}_3 = \frac{1}{\sqrt{m}} \begin{Bmatrix} -0.135 \\ -0.821 \\ 0.083 \end{Bmatrix}$$

Solve forced vibration problem $\bar{x} = \bar{X} e^{i\Omega t}$

$$(-\Omega^2 [m] + i\Omega [c] + [k]) e^{i\Omega t} = F_0 e^{i\Omega t}$$

Recast as

$$-\Omega^2 m [\bar{m}] + i\Omega c [\bar{c}] + k [\bar{k}] = \bar{F}_0, \text{ where}$$

$$[\bar{m}] = \begin{bmatrix} 7/2 & -1/2 & 0 \\ -1/2 & 3/2 & 0 \\ 0 & 0 & 5 \end{bmatrix}, \quad [\bar{c}] = \begin{bmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 1 \end{bmatrix},$$

$$[\bar{k}] = \begin{bmatrix} 1 & 0 & 0 \\ 0 & +3 & -1 \\ 0 & -1 & 1 \end{bmatrix}$$

divide the matrix equation by

$$\left(-\Omega^2 m/k [\bar{m}] + i\Omega c/k [\bar{c}] + [\bar{k}] \right) \bar{X} = \bar{F}$$

$$\left(\underbrace{-\Omega^2 m/k}_{\omega^2} [\bar{m}] + i \underbrace{\Omega c/k}_{3\mu} \frac{\sqrt{m}}{\sqrt{m}} \frac{1}{\sqrt{k}} \frac{1}{\sqrt{k}} [\bar{c}] + [\bar{k}] \right) \bar{X} = \bar{F}$$

$$\left(-\omega^2 [\bar{m}] + i 3\mu [\bar{c}] + [\bar{k}] \right) \bar{X} = \bar{F}$$

$$\begin{bmatrix} -\frac{7}{2}M^2 + 9\mu + 1 & \frac{1}{2}M^2 & 0 \\ \frac{1}{2}M^2 & -\frac{3}{2}M^2 + 3 & -3\mu \\ 0 & -3\mu & 5M^2 + 9\mu + 1 \end{bmatrix} \begin{Bmatrix} X_{1p} \\ X_{2p} \\ X_{3p} \end{Bmatrix} =$$

$$\begin{Bmatrix} 0 \\ 0 \\ F_0 \end{Bmatrix}$$

$$X_{1p} = \frac{\begin{vmatrix} 0 & \frac{1}{2}M^2 & 0 \\ 0 & -\frac{3}{2}M^2 + 3 & -3\mu \\ F_0 & -3\mu & 5M^2 + 9\mu + 1 \end{vmatrix}}{\Delta} = \frac{-F_0 M^2}{\Delta}$$

$$X_{2p} = \frac{\begin{vmatrix} -\frac{7}{2}M^2 + 9\mu + 1 & 0 & 0 \\ \frac{1}{2}M^2 & 0 & -3\mu \\ 0 & F_0 & 5M^2 + 9\mu + 1 \end{vmatrix}}{\Delta} = \frac{F_0(-M^2 + 25\mu + 2)}{2\Delta}$$

$$X_{3p} = \frac{\begin{vmatrix} -\frac{7}{2}M^2 + 9\mu + 1 & \frac{1}{2}M^2 & 0 \\ \frac{1}{2}M^2 & -\frac{3}{2}M^2 + 3 & 0 \\ 0 & -3\mu & F_0 \end{vmatrix}}{\Delta} = \frac{F_0((10M^4 - 24\mu^2 + 6) + (63\mu - 35M^2)\mu)}{2\Delta}$$

$$\Delta = (-25M^6 + \frac{3}{2}5^2M^4 + 65M^4 - 3M^25^2 - 4\frac{1}{2}M^2 + 2) + (25\frac{1}{2}3M^5 + 57\frac{1}{2}3M^3 + 53\mu)\mu$$

Resonance occurs at $\Omega = \omega_1, \omega_2$ and ω_3 . In nondimensional terms . . .

$$u_1 = \omega_1 / \sqrt{k/m} = 0,354$$

$$u_2 = \omega_2 / \sqrt{k/m} = 0,533$$

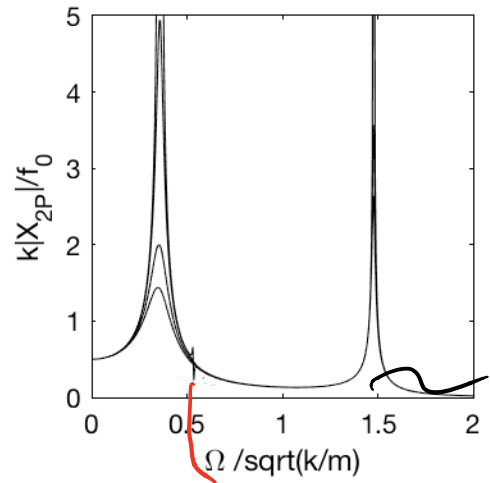
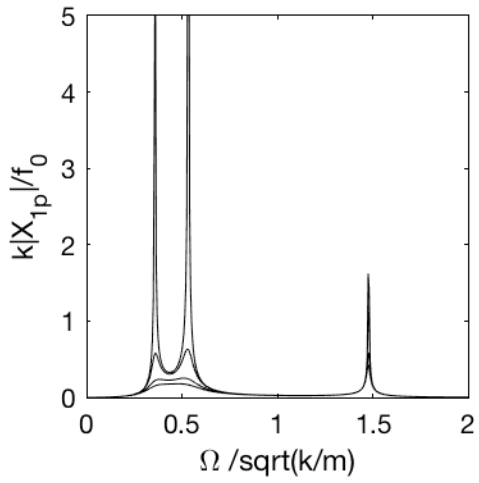
$$u_3 = \omega_3 / \sqrt{k/m} = 1,479$$

Antiresonance occurs at $\xi = 0$, when the response goes to zero

$$A_{1D} = \frac{-F_0 m R}{2\Delta}$$

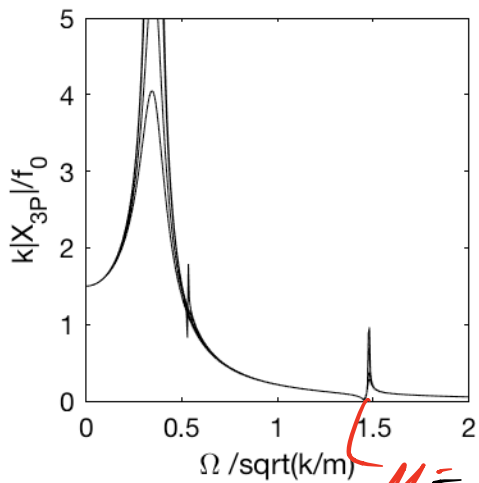
$$A_{2D} = \frac{F_0(-7m^2 + 2)}{2\Delta} \Rightarrow u = \sqrt{24} = 0,539$$

$$A_{3D} = \frac{F_0(110m^4 + 6 - 24m^2)}{2\Delta} \Rightarrow u = 0,532, 1,45$$



$\mu = 1.45$

sharp notch
 $\mu = 0.534$



$\mu = 1.45$

```

close all
clear all
clc

fprintf(['\n\nStarting file >>' mfilename '<< at ' datestr(now,0) '\n\n']);
format long
strp = [0 0 1; 0.5 0.5 0.8; 0 0.5 0.6; 0 0.5 0.4; 0 0.5 0.2];
syms m k omega mu zeta Fo

FS =12;

% Free Vibration Analysis
MM = [ 7/2 -1/2 0 ;
      -1/2 3/2 0 ;
      0 0 5 ];

CC = [ 1 0 0;
      0 0 0;
      0 0 1];

KK = [1 0 0 ;
      0 3 -1 ;
      0 -1 1 ];

[X,d] =eig(KK, MM);
[omegan,id] = sort(sqrt(diag(d)))
X= X(:,id)

% Symbolic analysis for CE
MMs = m*MM
KKs = k*KK

Mass =-mu^2*MM
Damp = zeta*i*mu*CC
Stiff = KK

DD = Mass+Damp+Stiff
Delta = simplify(det(DD))

%Cramer's Rule
DD1 = DD;
DD1(1,1) = 0; DD1(2,1) = 0; DD1(3,1)= Fo
Delta1 = simplify(det(DD1))

DD2 = DD;
DD2(1,2) = 0; DD2(2,2) = 0; DD2(3,2)= Fo
Delta2 = simplify(det(DD2))

DD3 = DD;
DD3(1,3) = 0; DD3(2,3) = 0; DD3(3,3)= Fo
Delta3 = simplify(det(DD3))

% Steady State Amplitudes
zeta_v = [0,0.2,0.5,0.7];
mu_v = linspace(0, 5, 1000)';

for ii = 1:length(zeta_v)

```

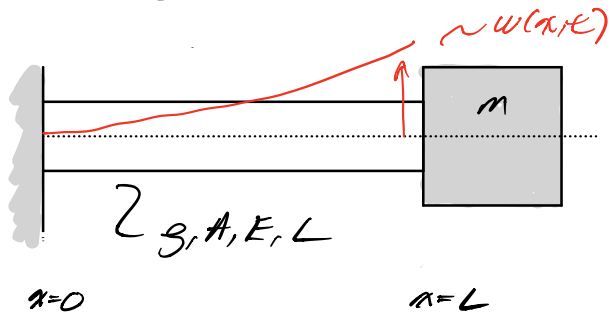
```
zeta = zeta_v(ii);
for jj = 1:length(mu_v)
    H = -mu_v(jj)^2*MM + i*zeta*mu_v(jj)*CC + KK;
    Y(:,jj) = inv(H)*[0;0;1];
end

figure(1)
subplot(1,3,1)
line(mu_v, abs(Y(1,:)), 'color', strp(ii,:)),hold on
axis([0,2,0,5]),xlabel('\omega /sqrt(k/m)'),ylabel('k|Y_1|/f_0')
box on
axis square

subplot(1,3,2)
line(mu_v, abs(Y(2,:)), 'color', strp(ii,:)),hold on
axis([0,2,0,5]),xlabel('\omega /sqrt(k/m)'),ylabel('k|Y_2|/f_0')
box on
axis square

subplot(1,3,3)
line(mu_v, abs(Y(3,:)), 'color', strp(ii,:)),hold on
axis([0,2,0,5]),xlabel('\omega /sqrt(k/m)'),ylabel('k|Y_3|/f_0')
box on
axis square
end
```


Problem ③



Note the governing equation

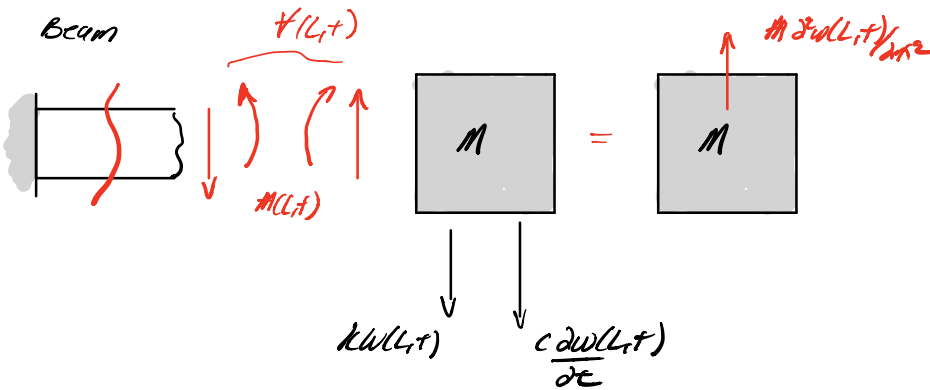
$$EI \frac{\partial^4 w}{\partial x^4} = -\rho A \frac{\partial^2 w}{\partial t^2}$$

The boundary conditions can be written as

at $x=0$

$$\begin{aligned} \text{① } w(0,t) &= \mathbb{I}(0) T(t) \rightarrow \mathbb{I}(0) = 0 \\ \text{② } \frac{\partial w}{\partial x}(0,t) &= \mathbb{I}'(0) T(t) \rightarrow \mathbb{I}'(0) = 0 \end{aligned}$$

at $x=L$, use ΣF and ΣM



$$\uparrow \Sigma M: M(L,t) = EI \frac{\partial^2 w(L,t)}{\partial x^2} = 0$$

$$\uparrow \Sigma F_y: V(L,t) - k w(L,t) - c \frac{\partial w(L,t)}{\partial t} = A \frac{\partial^2 w(L,t)}{\partial x^2}$$

$$\textcircled{3} \quad EI \frac{\partial^2 w(L,t)}{\partial x^2} = EI W''(L) T(t) = 0 \longrightarrow W''(L) = 0$$

$$\textcircled{4} \quad EI \frac{\partial^2 w(L,t)}{\partial x^2} - k w(L,t) - c \frac{\partial w(L,t)}{\partial t} = M \frac{\partial^2 w(L,t)}{\partial t^2}$$

assuming

$$T(t) = C e^{\gamma w t}$$

$$\dot{T}(t) = +\gamma w C e^{\gamma w t}$$

$$\ddot{T}(t) = -w^2 C e^{\gamma w t}$$

Then

$$\textcircled{4} \quad EI W'''(L) T(t) - k W(L) T(t) - c W(L) T'(t) = M W(L) T''(t)$$

$$EI W'''(L) T(t) - k W(L) T(t) - c \gamma w W(L) T(t) = -w^2 M W(L) T(t)$$

$$EI W'''(L) - k W(L) - \gamma c w W(L) = -w^2 M W(L)$$

The boundary conditions are ...

$$W(0) = 0$$

$$W'(0) = 0$$

$$W''(L) = 0$$

$$EI W'''(L) - k W(L) - \gamma c w W(L) = -w^2 M W(L)$$

Note the last equation is problematic, we can split between real and imaginary parts

$$R: EI W'''(L) - k W(L) + w^2 M W(L) = 0$$

$$I: \gamma c w W(L) = 0$$

The imaginary equation implies that

$W(L) = 0$, and from the real equation we

have

Now, the undamped condition is equivalent to looking at real part of the 4th boundary condition

$$EI W''''(L) - kW(L) + \omega^2 M W(L) = 0$$

OK,

$$W(x) = a \cosh Bx + b \sinh Bx + c \cos Bx + d \sin Bx$$

$$\text{where } B^2 = \sqrt{\frac{3A}{EI}} \omega$$

Thus the previous boundary conditions can be written as

$$\textcircled{1} W(0) = 0$$

$$\textcircled{2} W'(0) = 0$$

$$\textcircled{3} W''(L) = 0$$

$$\textcircled{4} EI W''''(L) - kW(L) + \omega^2 M W(L) = 0$$

$$\hookrightarrow W''''(L) = \frac{(k - \omega^2 M) W(L)}{EI}$$

Now let's apply

$$W(0) = a \cosh 0 + b \sinh 0 + c \cos 0 + d \sin 0 = 0$$

$$W'(0) = B(a \sinh 0 + b \cosh 0 - c \sin 0 + d \cos 0) = 0$$

$$0 = a + c \quad \rightarrow \quad c = -a$$

$$0 = B(b + d) \quad \rightarrow \quad d = -b \quad B \neq 0$$

$$W(x) = a (\cosh Bx - \cos Bx) + b (\sinh Bx - \sin Bx)$$

$$W'(x) = Ba (\sinh Bx + \sin Bx) + Bb (\cosh Bx - \cos Bx)$$

$$W''(x) = B^2 a (\cosh Bx + \cos Bx) + B^2 b (\sinh Bx + \sin Bx)$$

$$W''''(x) = B^3 a (\sinh Bx - \sin Bx) + B^3 b (\cosh Bx + \cos Bx)$$

$$\textcircled{3} \quad 0 = B^2 a (\cosh BL + \cos BL) + B^2 b (\sinh BL + \sin BL)$$

$$\textcircled{4} \quad 0 = B^3 a (\sinh BL - \sin BL) + B^3 b (\cosh BL + \cos BL) + \left(\frac{\omega^2 M - K}{EI} \right) (a (\cosh BL - \cos BL) + b (\sinh BL - \sin BL))$$

↳ look at this term

$$\text{Recall } \omega = \sqrt{\frac{EI}{\cancel{SA}}} B^2 \rightarrow \omega^2 = \frac{EI}{\cancel{SA}} B^4$$

then

$$\begin{aligned} \frac{\omega^2 M - K}{EI} &= \frac{\omega^2 M}{EI} - \frac{K}{EI} = \frac{EI}{\cancel{SA}} \frac{\cancel{AA}}{\cancel{EI}} B^4 - \frac{K}{EI} \\ &= \frac{AA}{\cancel{SA}} B^4 - \frac{K}{EI} \end{aligned}$$

$$\textcircled{4} \quad 0 = B^3 a (\sinh BL - \sin BL) + B^3 b (\cosh BL + \cos BL) + \left(\frac{AA B^4}{\cancel{SA}} - \frac{K}{EI} \right) (a (\cosh BL - \cos BL) + b (\sinh BL - \sin BL))$$

$$0 = (BL)^3 a (\sinh BL - \sin BL) + (BL)^3 b (\cosh BL + \cos BL) +$$

$$\left(\frac{AA (BL)^4}{\cancel{SA}} - \frac{KL^3}{EI} \right) (a (\cosh BL - \cos BL) + b (\sinh BL - \sin BL))$$

$$\text{Recall } AA = \cancel{SA} L \text{ and } \alpha = \frac{KL^3}{EI}$$

$$0 = (BL)^3 a (\sinh BL - \sin BL) + (BL)^3 b (\cosh BL + \cos BL) + (BL)^4 - \alpha (a (\cosh BL - \cos BL) + b (\sinh BL - \sin BL))$$

Write as a matrix ...

$$\begin{bmatrix} \cosh TBL + \cos TBL & \sinh TBL + \sin TBL \\ (TBL)^3 (\sinh TBL - \sin TBL) + (TBL)^4 - a (\cosh TBL - \cos TBL) & (TBL)^3 (\cosh TBL + \cos TBL) + (TBL)^4 - a (\sinh TBL - \sin TBL) \end{bmatrix} \begin{Bmatrix} a \\ b \end{Bmatrix} = \vec{0}$$


$$D$$

$$\Delta = \det(D) = CE$$

Now, solve for modeshapes using 1st row of matrix equation

$$(\cosh BL + \cos BL)a + (\sinh BL + \sin BL)b = 0$$

$$b = \frac{(\cosh BL + \cos BL)a}{(\sinh BL + \sin BL)}$$

$$W(x) = a(\cosh Bx - \cos Bx) + b(\sinh Bx - \sin Bx)$$

$$W(x) = a(\cosh Bx - \cos Bx) + \frac{(\cosh BL + \cos BL)a}{(\sinh BL + \sin BL)}(\sinh Bx - \sin Bx)$$

If we consider damping the $cw W(L) = 0$

this implies $W(L) = 0$

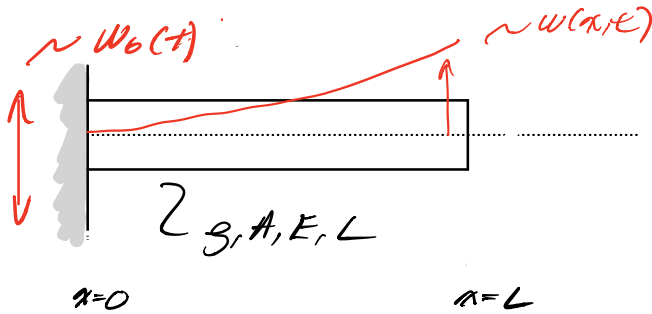
$$W(L) = a(\cosh BL - \cos BL) + \frac{(\cosh BL + \cos BL)a}{\sinh BL + \sin BL}(\sinh BL - \sin BL) = 0$$

This implies $a = 0$

∴ damping at boundaries doesn't affect boundary conditions since the result would be no motion

See code

Consider,



Note the governing equation $EI \frac{\partial^4 w}{\partial x^4} = -\rho A \frac{\partial^2 w}{\partial t^2}$

Now, $w(x,t) = w_r(x,t) + w_0(t)$

$$\frac{\partial^4 w}{\partial x^4} = \frac{\partial^4 w_r}{\partial x^4}$$

$$\frac{\partial^2 w}{\partial t^2} = \frac{\partial^2 w_r}{\partial t^2} + \ddot{w}_0(t)$$

$$EI \frac{\partial^4 w_r}{\partial x^4} = -\rho A \frac{\partial^2 w_r}{\partial t^2} - \rho A \ddot{w}_0(t)$$

$$EI \frac{\partial^4 w_r}{\partial x^4} + \rho A \ddot{w}_0(t) = \rho A \frac{\partial^2 w_r}{\partial t^2}$$

The term $\rho A \ddot{w}_0(t)$ acts like distributed forcing and does not change the boundary conditions.

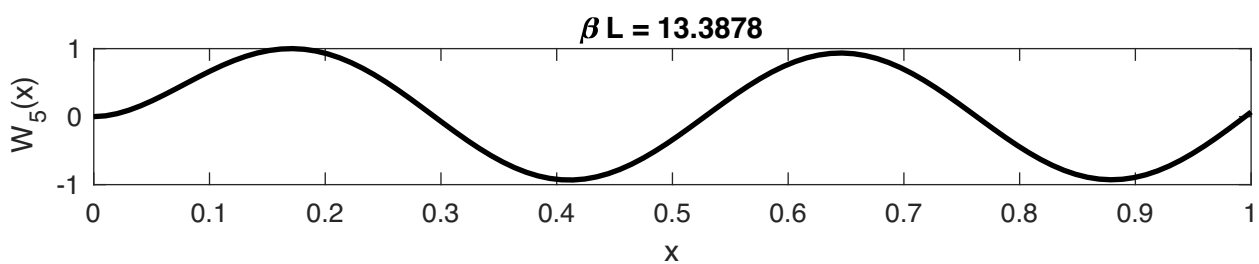
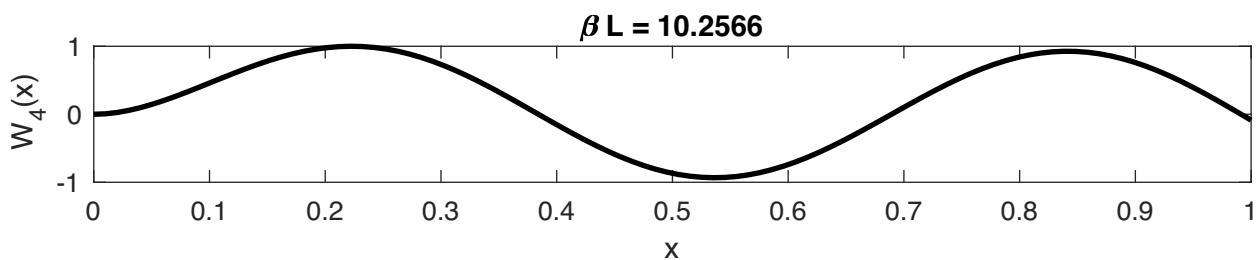
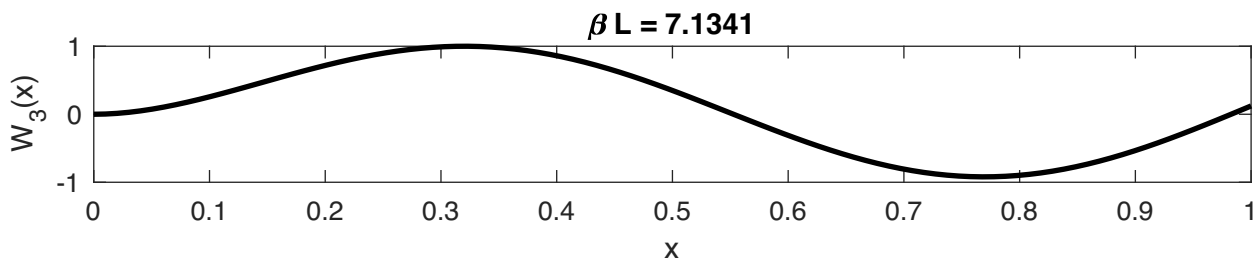
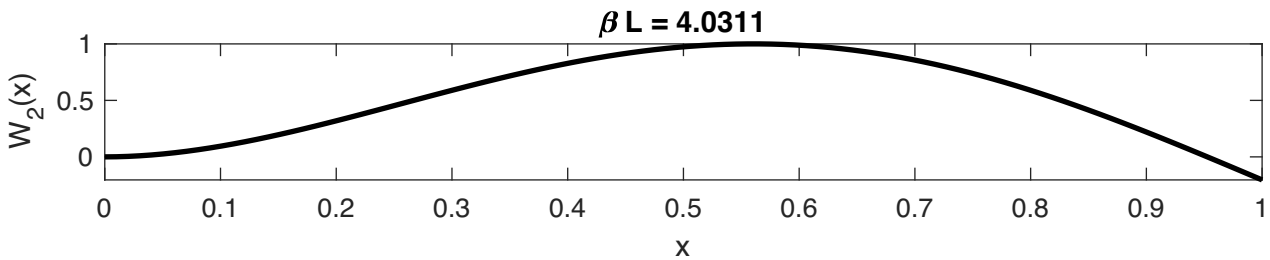
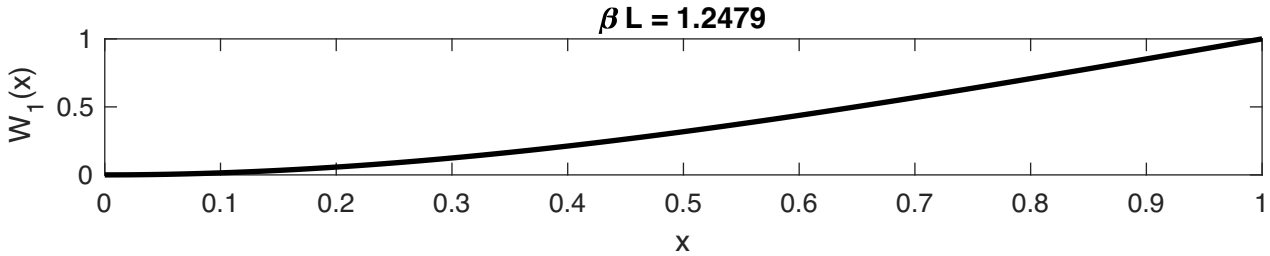
The boundary conditions can be written as

$$w_r(0,t) = 0, \quad EI \frac{\partial^2 w_r}{\partial x^2} = 0$$

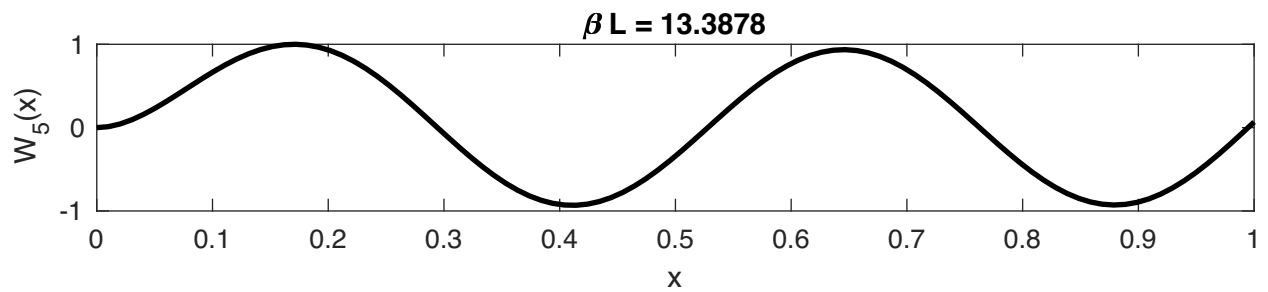
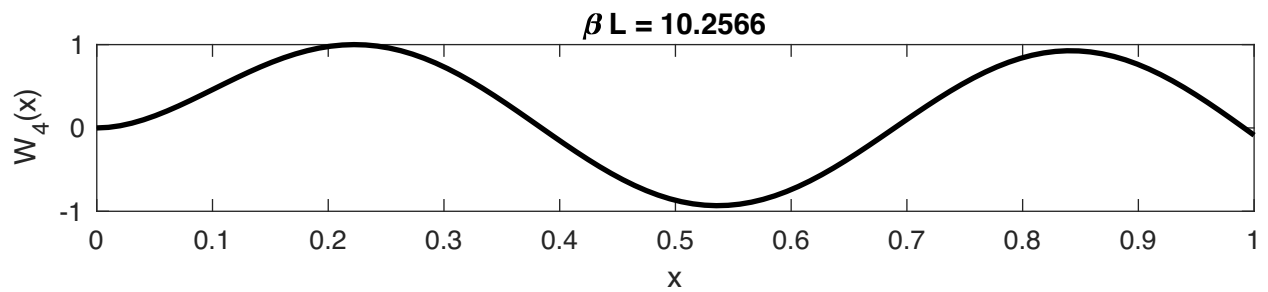
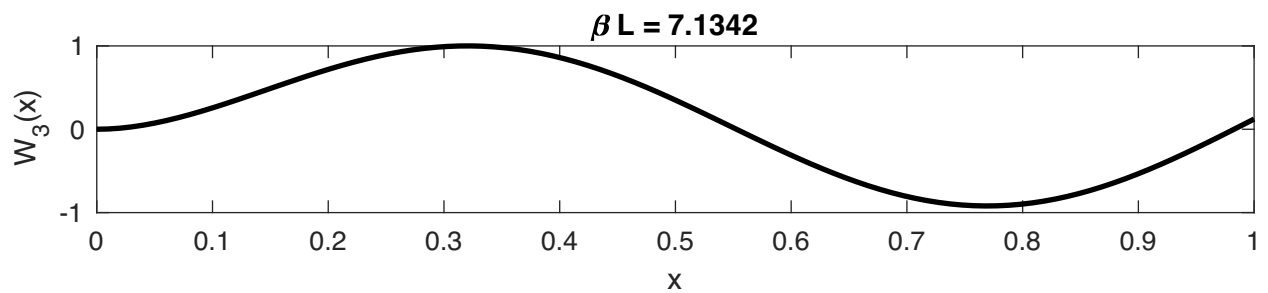
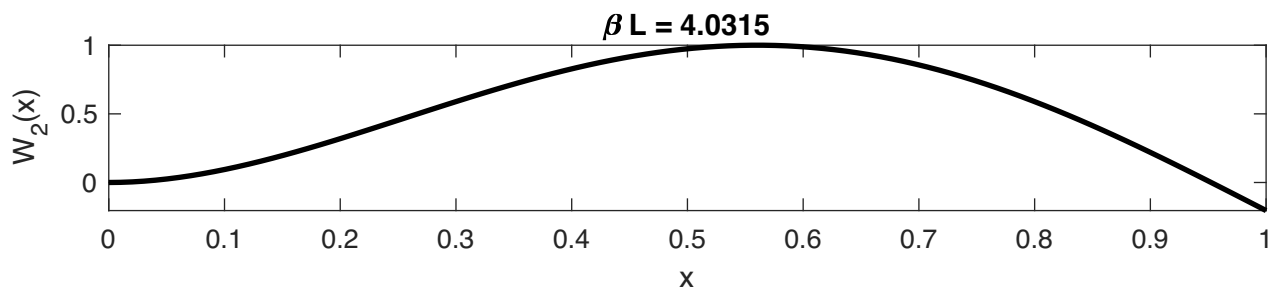
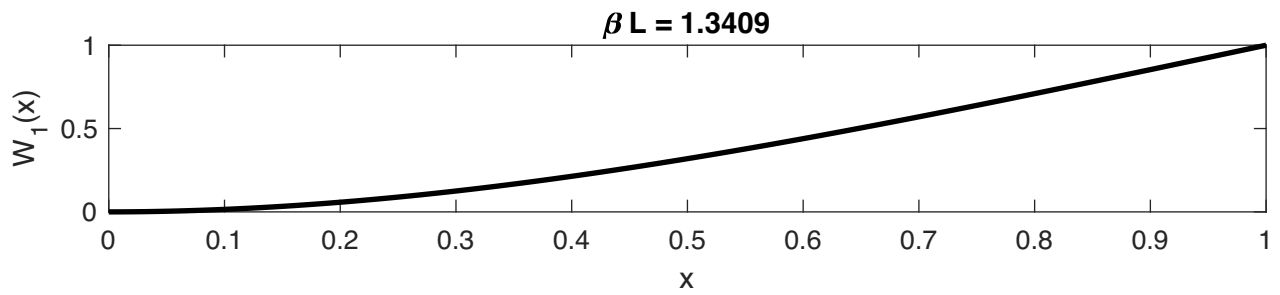
$$\dot{w}_r(0,t) = 0, \quad EI \frac{\partial^3 w_r}{\partial x^3} = 0$$

same as a fixed free beam.

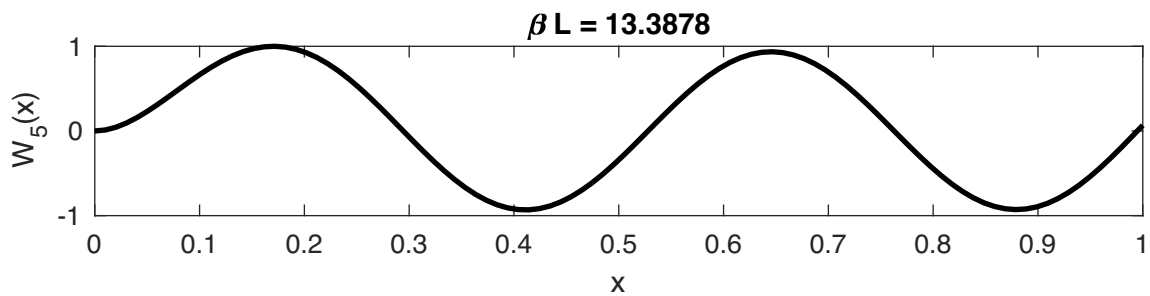
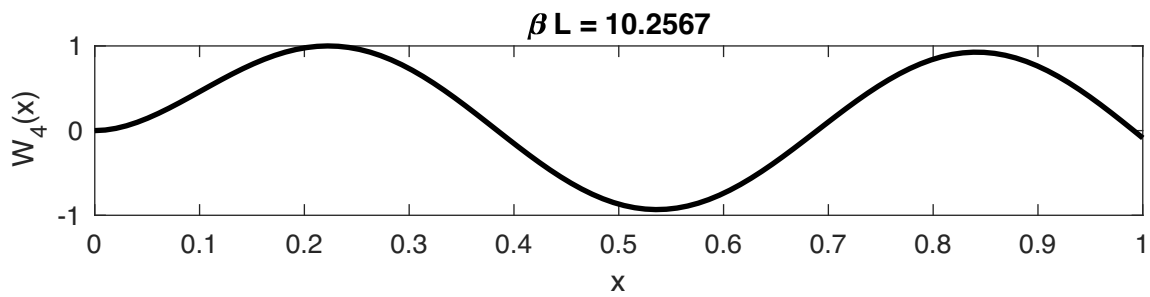
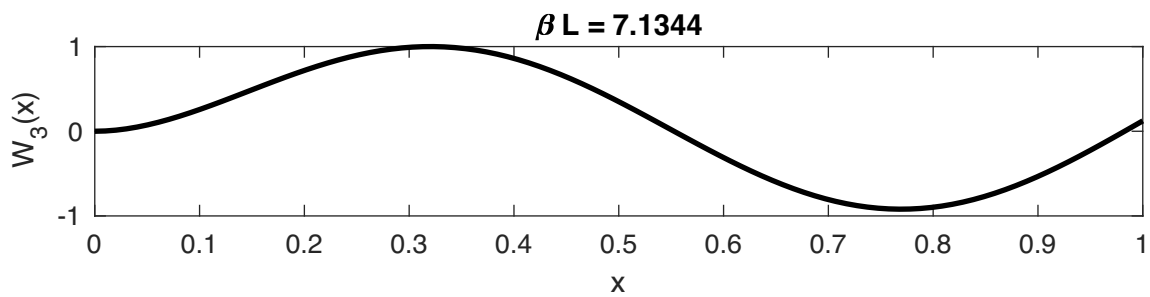
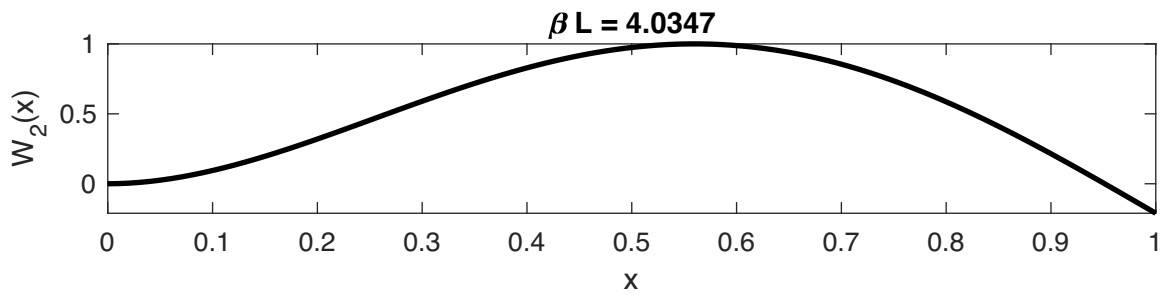
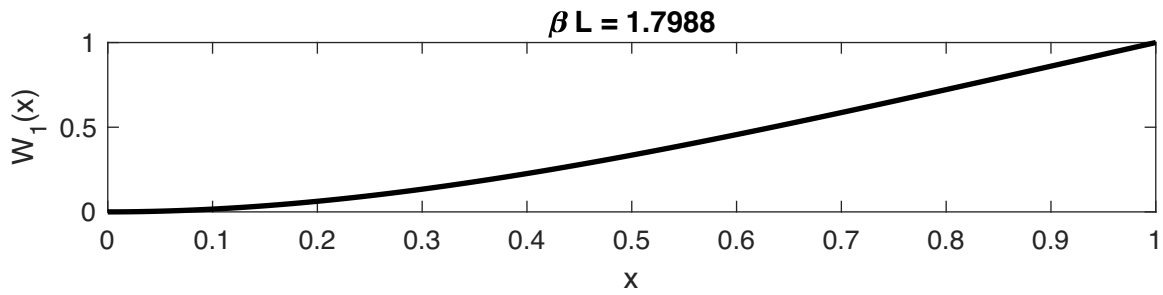
$\alpha = 0$ same as HW 5.3



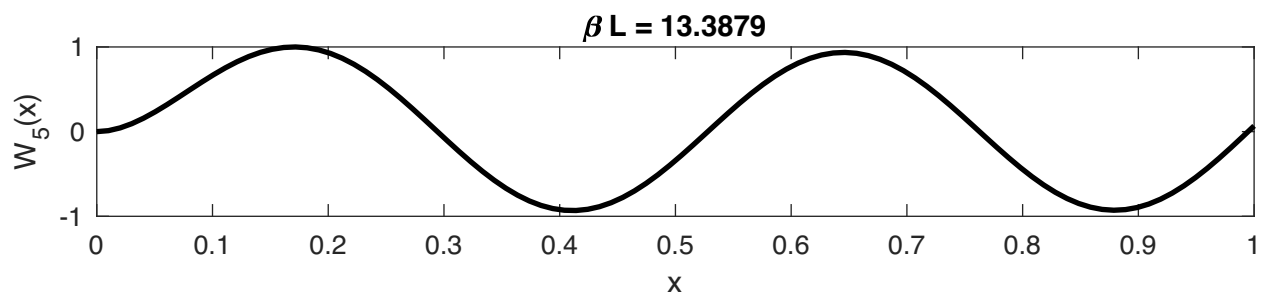
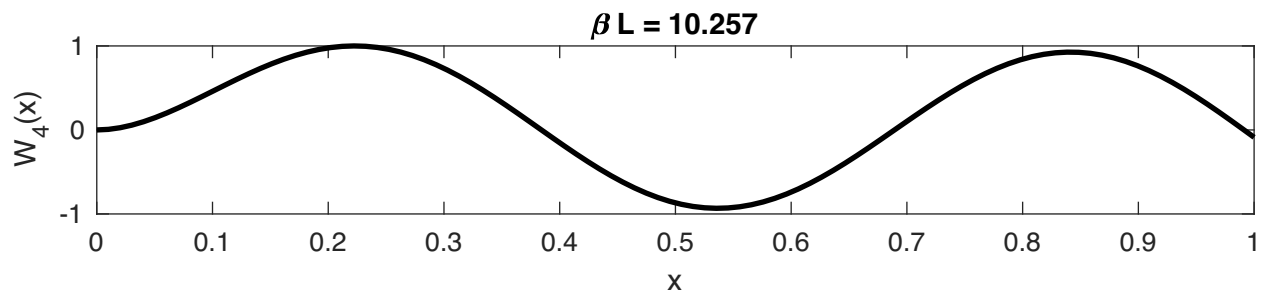
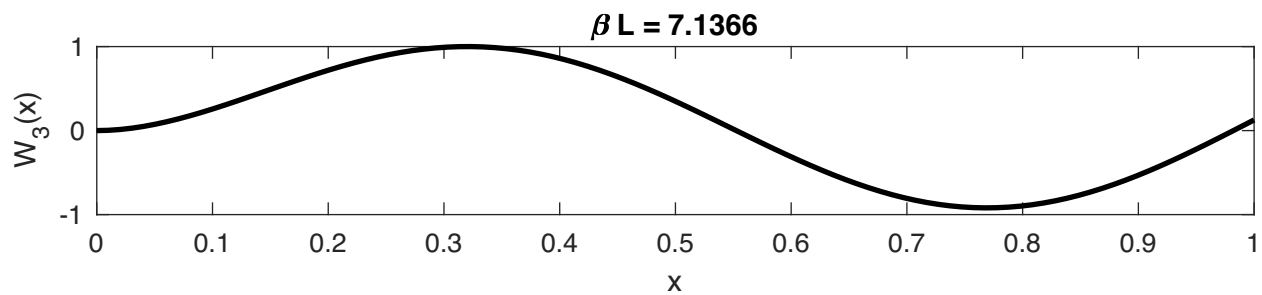
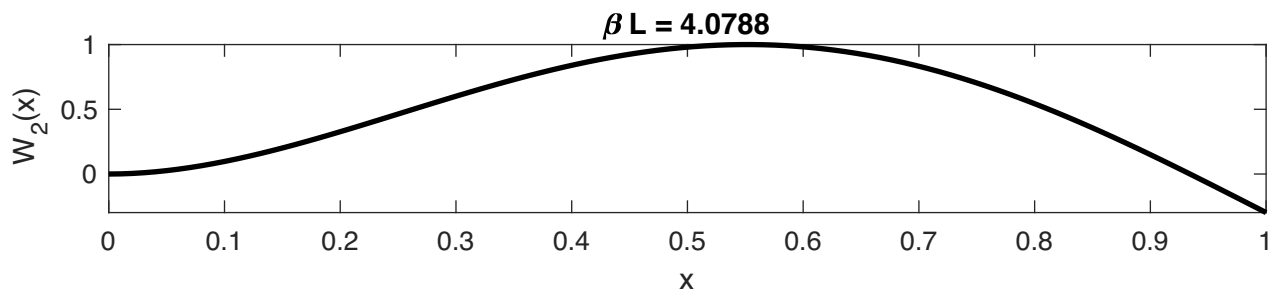
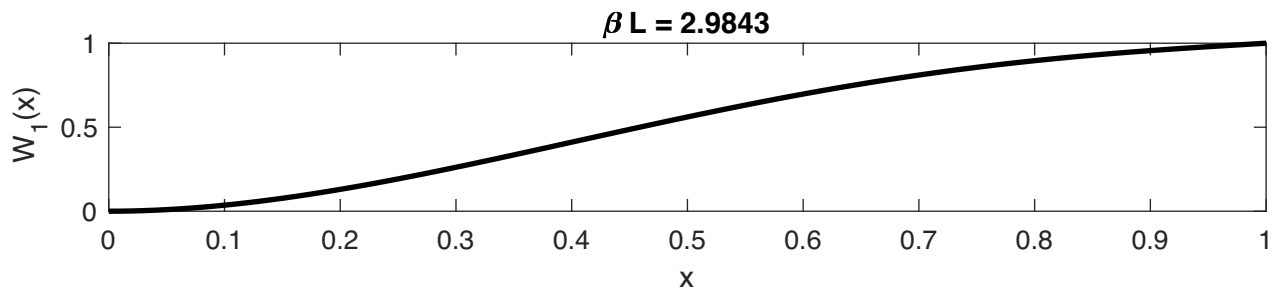
$$\alpha = 1$$



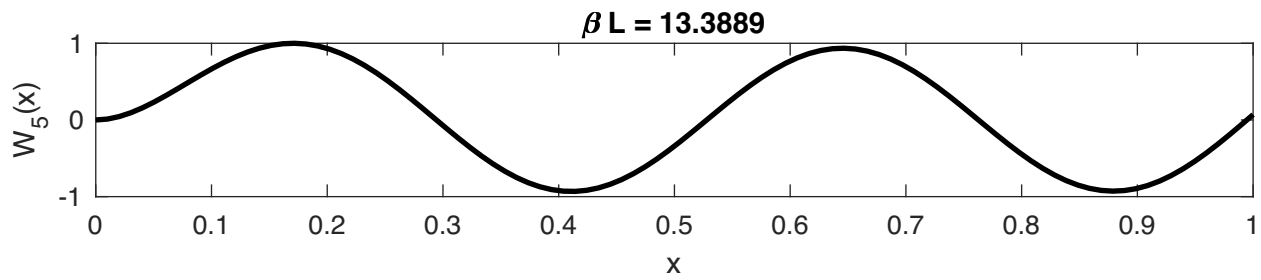
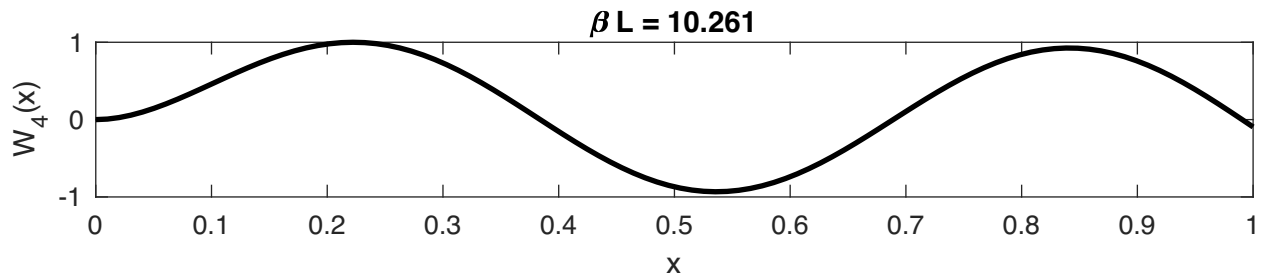
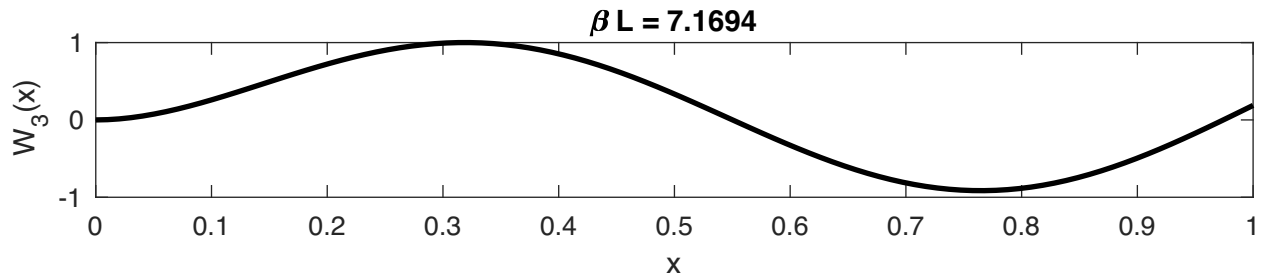
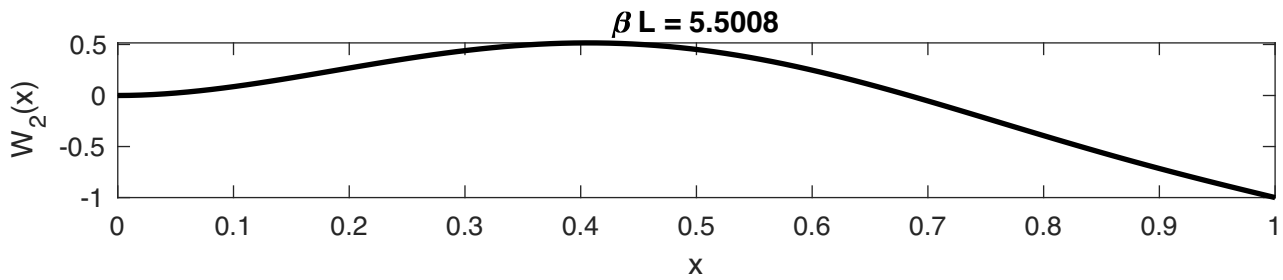
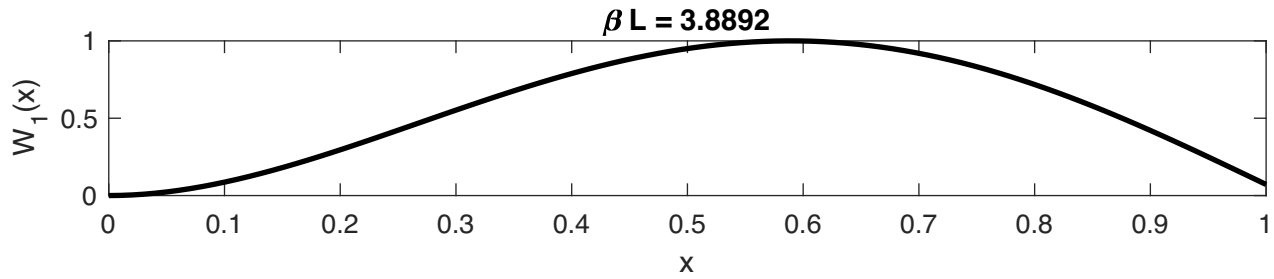
$$a = 10$$



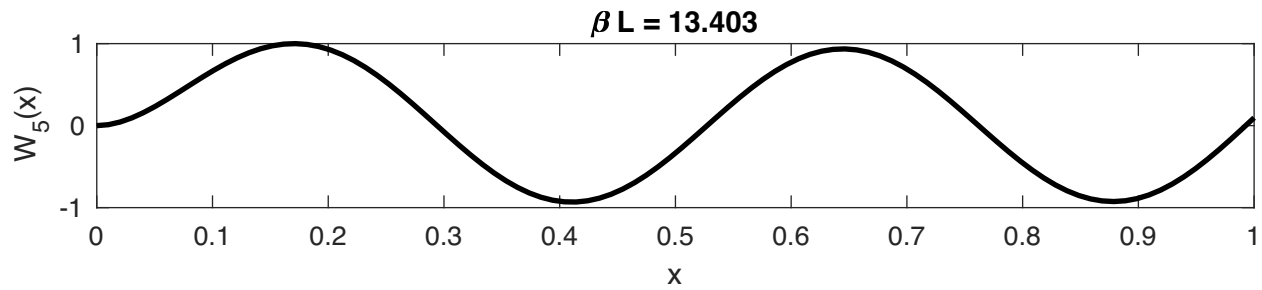
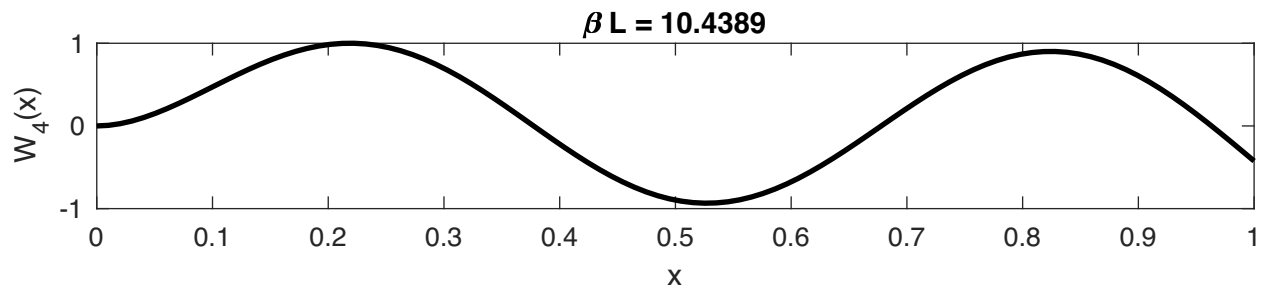
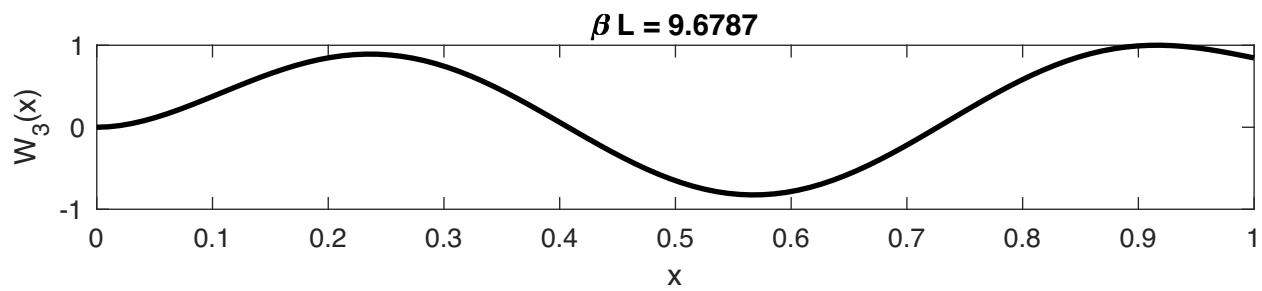
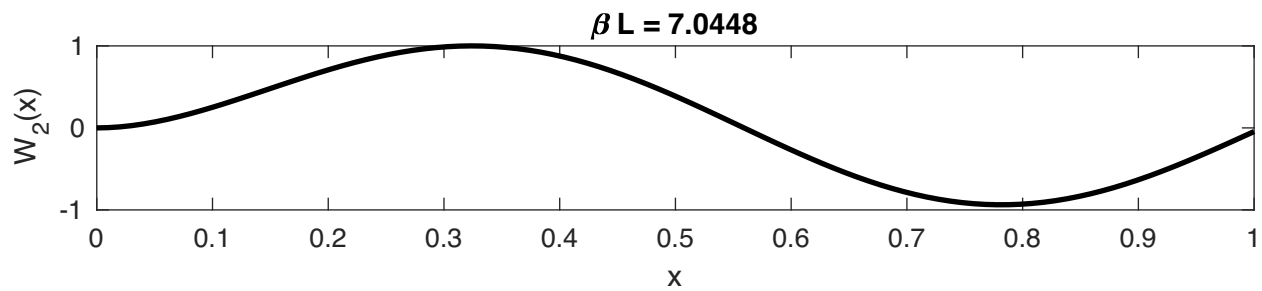
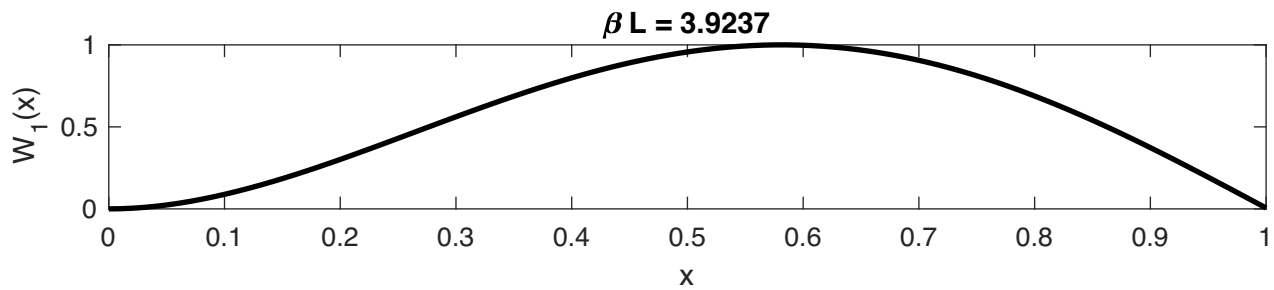
$$\alpha = 100$$



$$\alpha = 1000$$



$$\alpha = 10000$$



```

clc
clear
close all
fprintf(['\n\n\nStarting file >>' mfilename '<< at ' datestr(now,0) '\n\n']);
format long
% Exam 2 Problem 3

% x= Beta*L
syms B x l a b c d M E I w rho L A BL alpha

W = a*(cosh(B*x)-cos(B*x))+b*(sinh(B*x)-sin(B*x))

% Write as matrix pull out coefficients
BC11 = cosh(BL)+cos(BL);
BC12 = sinh(BL)+sin(BL);
BC21 = (BL).^3.*(sinh(BL)-sin(BL))+((BL).^4 - alpha).*(cosh(BL)-cos(BL));
BC22 = (BL).^3.*(cosh(BL)+cos(BL))+((BL).^4 - alpha).*(sinh(BL)-sin(BL));

BC = [BC11 BC12; BC21 BC22];

CE = simplify(det(BC));
CE = matlabFunction(CE);
%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%
alpha = 0;
CEfun = @(BL) CE(BL,alpha); % This is for optimization
%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%
BL = linspace(0,5*pi,10^4);
CEv = CE(BL,alpha);

% This CE does not have plot that we can easily identify our initial
% guess
% In order to get initial guess we look at places where there is a sign
% change in the CE
[Val1, loc1] = find(abs(diff(sign(CEv)))==2); % This time loc1 is indices
%
figure
plot(BL(1:end),sign(CEv),'r')
axis([0 inf -1.2 1.2])

% Initial Guesses
x0v = BL(loc1);
options_all=[];
% set your length L
L=1;
for i =1:5
    x0 = x0v(i);
    fprintf(['\n\n\n >> Mode' num2str(i) '<< \n\n']);
    % Fzero
    [xval, fval, exitflag] = fzero(CEfun, x0, options_all);
    betaL(i) = xval;%This is Beta*L
    beta(i) = betaL(i)/L; % This is Beta
end

%Solve for Modeshapes
%solve for a constant
bsol = -BC11/BC12;
%plug into W(x)
Wmode= simplify(subs(W,b,bsol));

```

```
pretty(Wmode)
%Plot mode shapes
x = linspace(0,L,100);
figure
for i =1:5
    B = beta(i);
    BL = betaL(i);
    a = 1;
    WmodeP = eval(Wmode);
    Wmax = max(abs(WmodeP));
    WmodeP = WmodeP/Wmax;
    subplot(5,1,i)
    line(x,WmodeP,'linewidth',2,'color','k')
    xlabel('x')
    ylabel(['W_',num2str(i),'(x)'])
    title(['\beta L = ', num2str(BL)])
    box on
end
```